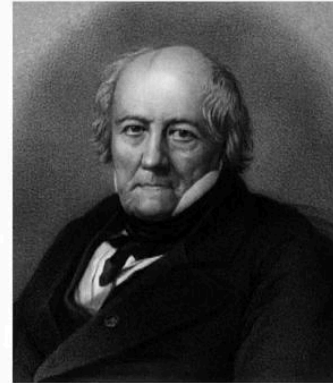


Thermodynamique

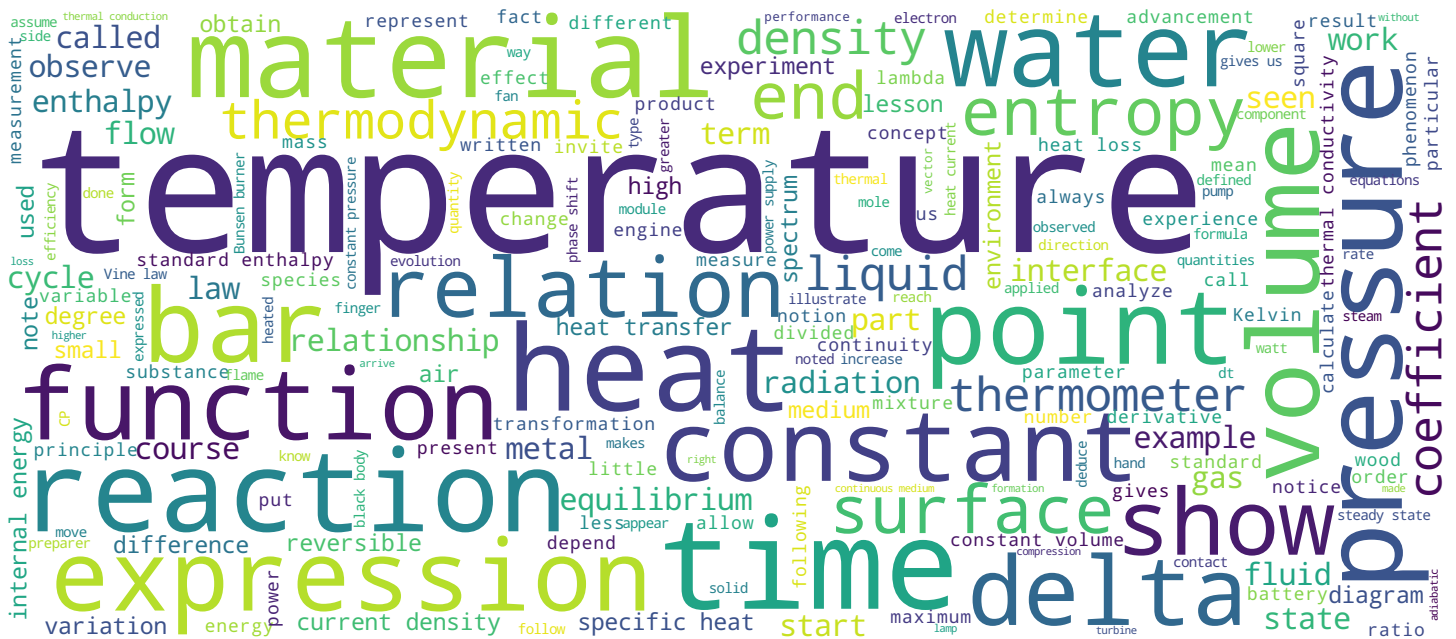
Expériences : Transferts thermiques



Jean-Baptiste Biot, 1774-1862



Prof. Jean-Philippe Ansermet



Video





- Conduction thermique
- Refroidissement avec ou sans ventilation
- Déphasage thermique
- Effusivité
- Rayonnement

Thermodynamique

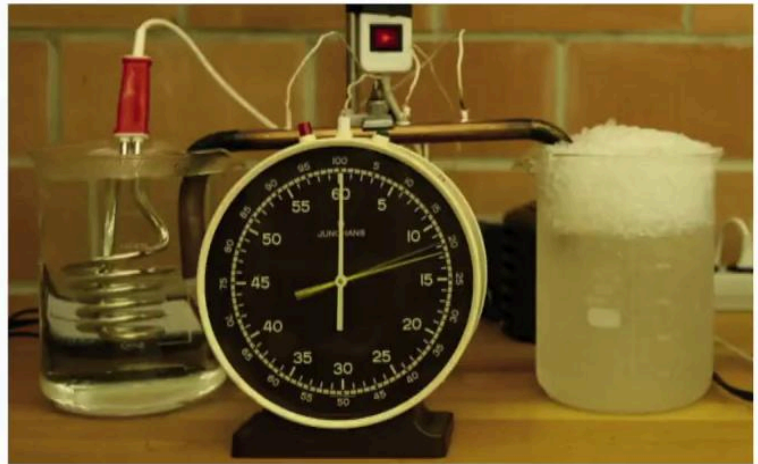
Here I am again to present you some experiences. Chantal Matteau and Martin Brochu give you a lesson on heat transfers in stationary and transient regimes. Here. I would like to illustrate their lesson by first showing them the conduction in a copper rod in steady state. Next, I would like to address the issue of heat loss of a metal in contact with air. Next, I would like to illustrate the concept of thermal phase shift. And then I would like to show you an experiment for which the notion of diffusivity intervenes. And I will end with two experiences on heat transfer by radiation starting with thermal conduction.

Notes

Summary



0m 04s



Thermodynamique

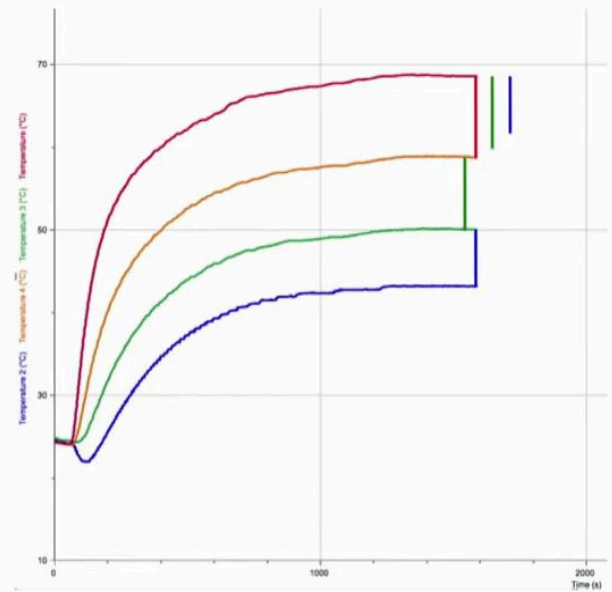
I suggest you look at the temperature along a copper rod. That you can see right behind that big timer. Four thermocouples were mounted at equal distances along the bar. We will record these temperatures while the bar is held at one end in boiling water and at the other end in crushed ice. Melting. So at zero degrees.

Notes

Summary



1m 03s



Thermodynamique

I show you on this graph. The result of the experiment. You have. On the ordinate the temperature in degrees Celsius. On the x-axis the time in seconds. We are not interested here in to the temporal evolution, but you see that we have after 1000 seconds, a stationary regime is reached. What I observe is that the temperature differences do not are not exactly what I had planned for the ideal case. If you want, I can plot the difference between two reported temperatures here. These gaps? You can see that with this little montage graph, we don't always have the same gap of temperature while the four thermometers are acquired distant.

Notes

Summary



1m 34s

$$j_Q = -\kappa \nabla T \quad P_Q = j_Q (\pi r^2)$$

Entre deux thermomètres :

α : perte relative entre deux thermomètre

S : surface de cuivre

Entre le cuivre et l'air : ΔT

$$h = \frac{\alpha P_Q}{\Delta T S} \text{ (W m}^{-2}\text{K}^{-1}\text{)}$$

$$\alpha \approx 15\%$$

Thermodynamique

How to analyze this experience? So I'm talking about Fourier's law. I don't call it thermal conductivity. If there is no heat loss, this law can be applied everywhere. We have the same thermal power P_Q . Which depends on the heat current, the heat current density G_Q and the cross-sectional area of the bar, then r square which gives us the power. We would have the same power everywhere. Obviously, this is not the case. Otherwise we would have the same ΔT between every three thermometers, every four thermometers. So I'm going to assume that in two thermometers, I have a relative loss of power, but I will call this relative loss α . I call the surface of the copper between two thermometers S . I'm going to guess, that's the parameter I'm least comfortable with. I'm going to assume that I have a difference of temperature ΔT between the surface of the metal and the air in its vicinity. And so I'll calculate a heat loss coefficient like this. Per unit area. I have a loss in watts per square meter per Kelvin. If I take my digital data, I have my little geometric construction that reappears here. Roughly speaking, I have a loss of about 15% between two thermometers.

Notes

Summary



2m 29s



$$j_Q = -\kappa \nabla T \quad P_Q = j_Q (\pi r^2)$$

Entre deux thermomètres :

α : perte relative entre deux thermomètre

S : surface de cuivre

Entre le cuivre et l'air : ΔT

$$h = \frac{\alpha P_Q}{\Delta T S} \text{ (W m}^{-2}\text{K}^{-1}\text{)}$$

$$\alpha \approx 15\% \quad \Delta T \approx 25 \text{ K}$$

$$\kappa = 380 \text{ W/(mK)} \quad \nabla T \approx \frac{98 \text{ K}}{0.35 \text{ m}}$$

$$r \approx 10 \text{ mm} \quad \Delta \ell \approx 0.05 \text{ m}$$

$$h \approx 64 \text{ W m}^{-2}\text{K}^{-1}$$

Thermodynamique

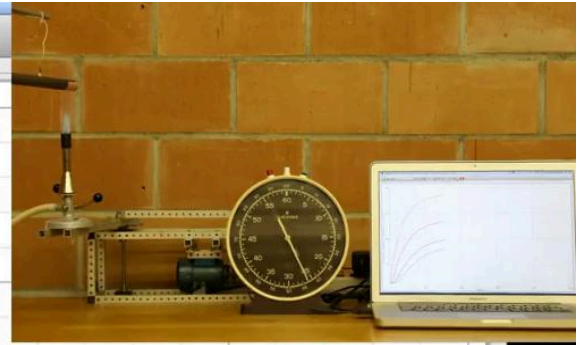
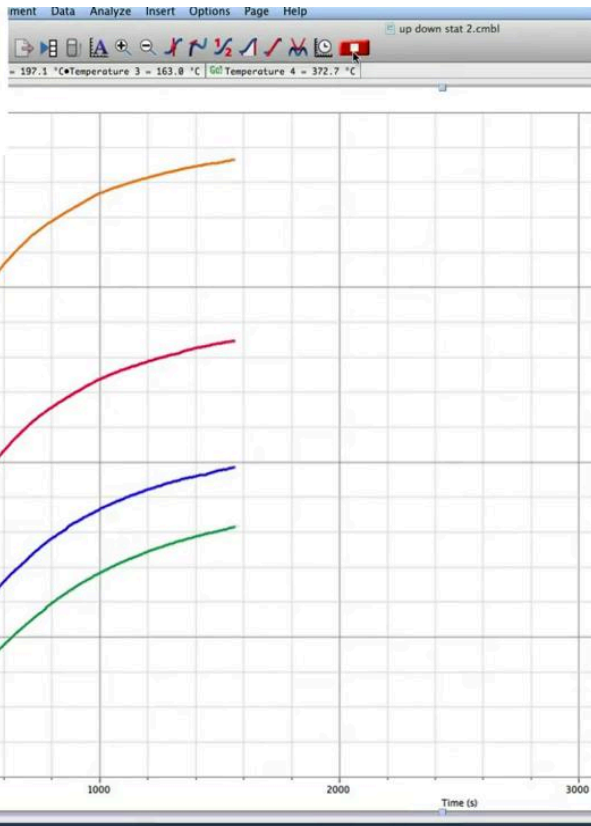
Of course, the loss is continuous along the bar. Here, I simply try to estimate this coefficient h . So look for an order of magnitude. So I simplify things for myself assuming there is a loss, just one location per segment. The Delta t . As I said earlier, I'm not sure of its value. I'll take 25, Kelvin. I will look up the thermal conductivity of copper in a table. I have water boiling at ninety-eight degrees. And then the length of the bar is about 35 cm. It gives me the temperature gradient. So I can deduce the thermal power. P_Q . I know the diameter of the tube. Finally, the copper bar. The distance between two thermometers is five cm. And so I deduce a loss of 64 watts per square meter per Kelvin. It may be a bit much compared to the tabular values, but at least we have an order of magnitude.

Notes

Summary



4m 11s



I am now passing another experience that you can analyze, that will allow you to better quantify this heat loss at the surface of a metal. We will consider a copper bar, again with four thermometers and this time we will observe. The evolution of the temperature. At the four measuring points. So the temperature as a function of time, with or without a fan that blows on the bar. Let's start with the static case. See here, you could see the Bunsen burner which heats the end of the bar and you have here now, as a function of time, the temperature of the four thermocouples. And you can see that we're getting close. A stationary state.

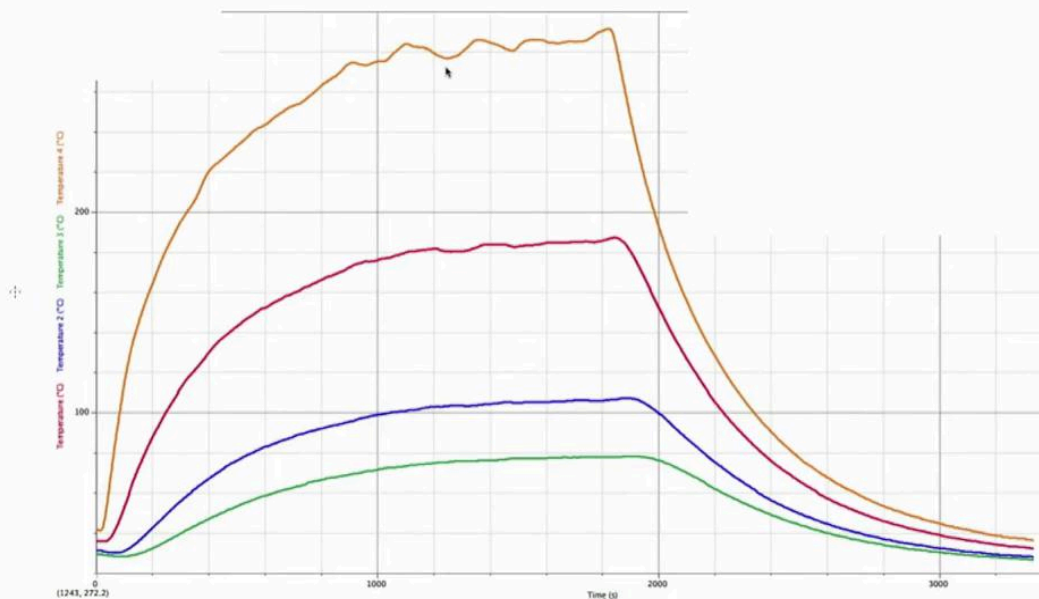
Notes

Summary



5m 20s

Refroidissement avec ventilation



Thermodynamique

I show you here. The whole of the measurements. Of course, halfway through, we stopped the flame. That's why the baron gets cold. And I invite you to use the equation of heat to account for these data. I would now like to see how these data change experimentally when the fan blows on the bar. Do you see the fan in operation here? And the temperature is measured again as a function of time. Mostly we observe, for the first thermometer closest to the flame, fluctuations that come from the fact that the fan disturbs the flame. I now show you the set of measures for this same bar. But with the fan on, you immediately notice that you reach a steady state, but temperatures are lower and I invite you to analyze the graph to observe that the speed cooling capacity is also greater.

Notes

Summary



6m 22s



I would now like to illustrate the concept of thermal phase shift. To do this, I propose to examine this brass bar once again. But this time, the Bunsen burner is mounted on an arm. Who also, what makes him periodically pass under the bar. The Bunsen burner that is far from the bar and here it is under the bar. And that he is once again stepping away from the bar.

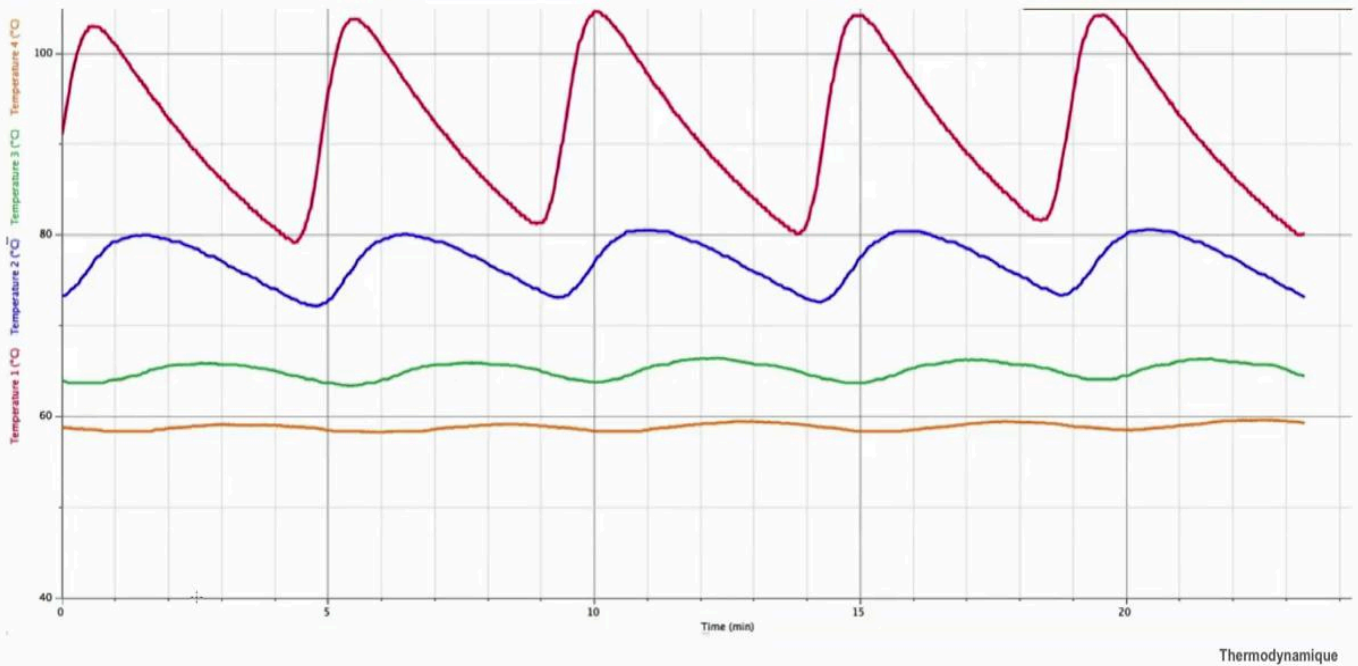
Notes

Summary



7m 36s

Déphasage thermique



I will show you the whole measure here. Note that the recording was made after several passes of the flame of the Bunsen burner under the bar to reach this periodic regime. What I invite to notice, is that obviously the highest temperature is that of the thermometer which is closest to the hot spring. And you notice that the further away we get of this hot spot, the more the maximum takes place by themselves. On the fourth curve, one can hardly distinguish an oscillation. But if you take the time to look at it carefully, you can see that there are always this phase shift which is more and more important. The further away from the hot spring.

Notes

Summary



8m 15s



Thermodynamique

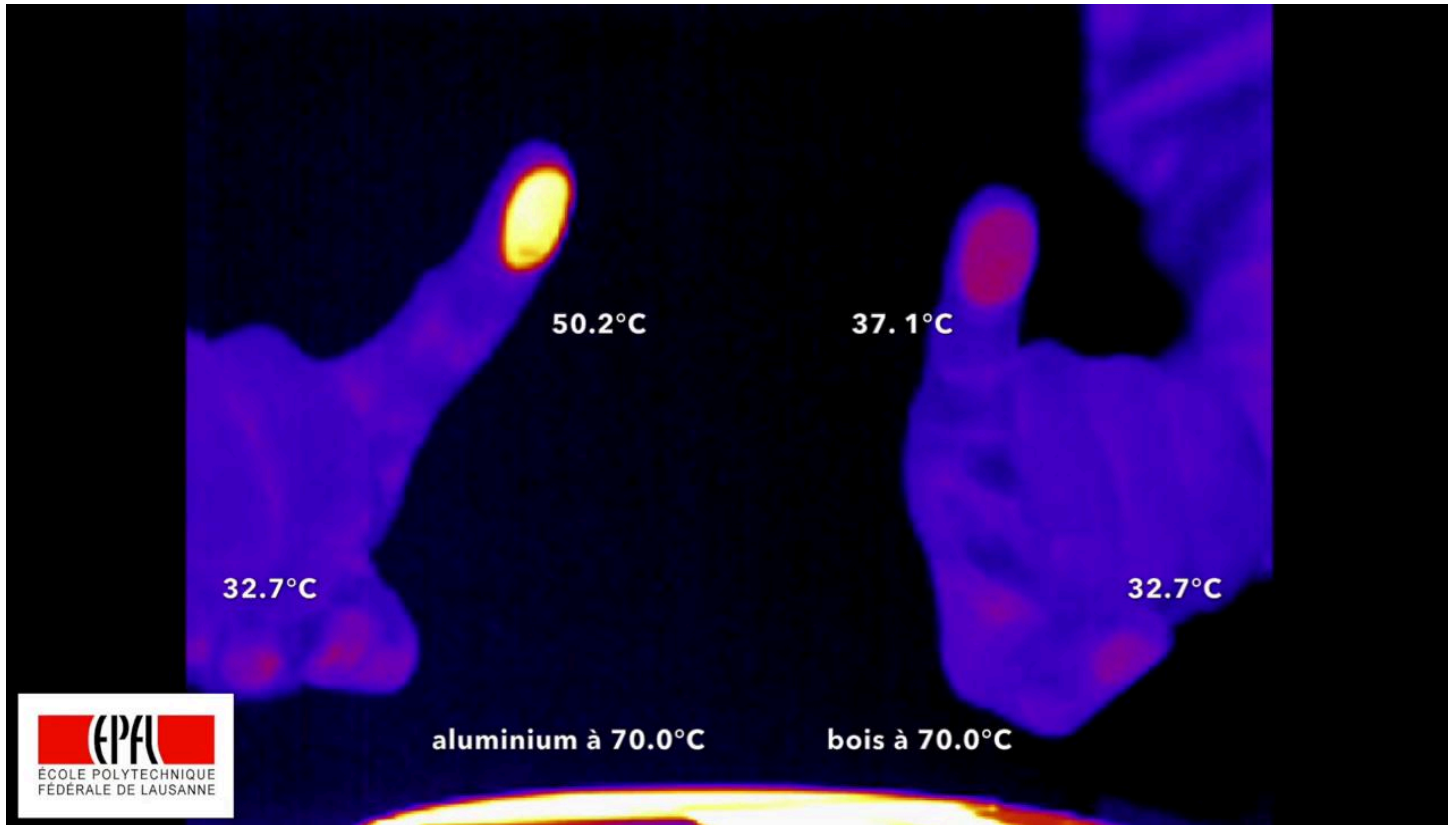
I would like to show you now an experiment for which the notion of diffusivity intervenes. Like this image. This image is an infrared camera image showing two fingers. It's two fingers that have been pressed on surfaces of different nature but which are at the same temperature.

Notes

Summary



9m 09s



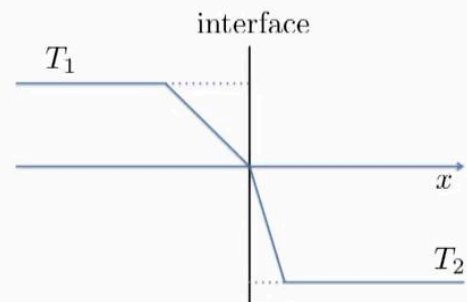
I invite you to watch the video. Here we see a block that is heated to seventy degrees Celsius on it. For a long time, a square of wood and a square of metal were arranged. The preparer will press his fingers on both materials. And then present them to the infrared camera. Here we see the detail. You notice that the temperature is decreasing, but slowly enough so that the initial measurement is representative of the temperature at the interface between the finger and the metal or wood. This temperature is used in the calibration of the camera. We see that on the wood, we have a much lower temperature than what we get on the metal.

Notes

Summary



9m 30s



$$T_i = \frac{E_1 T_1 + E_2 T_2}{E_1 + E_2}$$

$$E_1 = \sqrt{\kappa_1 c_1} \quad E_2 = \sqrt{\kappa_2 c_2}$$

$\kappa_i (i = 1, 2)$: conductivité thermique

$c_i (i = 1, 2)$: chaleur spécifique

Thermodynamique

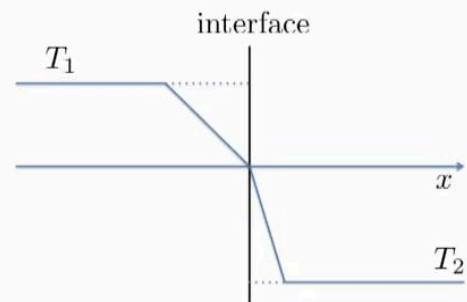
How do you do it? Describe the phenomenon. I propose the following description. I represent here the interface between the two materials. Initially, one material has a temperature T_1 , the other has a temperature T_2 , so T_1 . It is for example the wood and T_2 is the must very quickly. When these two materials are in contact with each other, they will not stay at different temperatures due to the effect of the of thermal conduction, even on a very thin layer around the interface. The interface will take a temperature. This is what I have represented in the drawing. There will of course be a gradient of temperature corresponding to a heat flow between the two materials. Now both materials are assumed to have different thermal conductivities. We have a continuity of the heat current at the interface, so we have slopes corresponding to different thermal conductivity. If we do the analysis from the heat equation, we will see that the parameter that intervenes is the flows. It is noted here them an index one or two which depends on the conductivity of the material and the specific heat of the material and a volumetric specific heat. The formula we get when we do short time analysis at a very small distance from the interface.

Notes

Summary



10m 27s



$$T_i = \frac{E_1 T_1 + E_2 T_2}{E_1 + E_2}$$

$$E_1 = \sqrt{\kappa_1 c_1} \quad E_2 = \sqrt{\kappa_2 c_2}$$

$\kappa_i (i = 1, 2)$: conductivité thermique

$c_i (i = 1, 2)$: chaleur spécifique

Thermodynamique

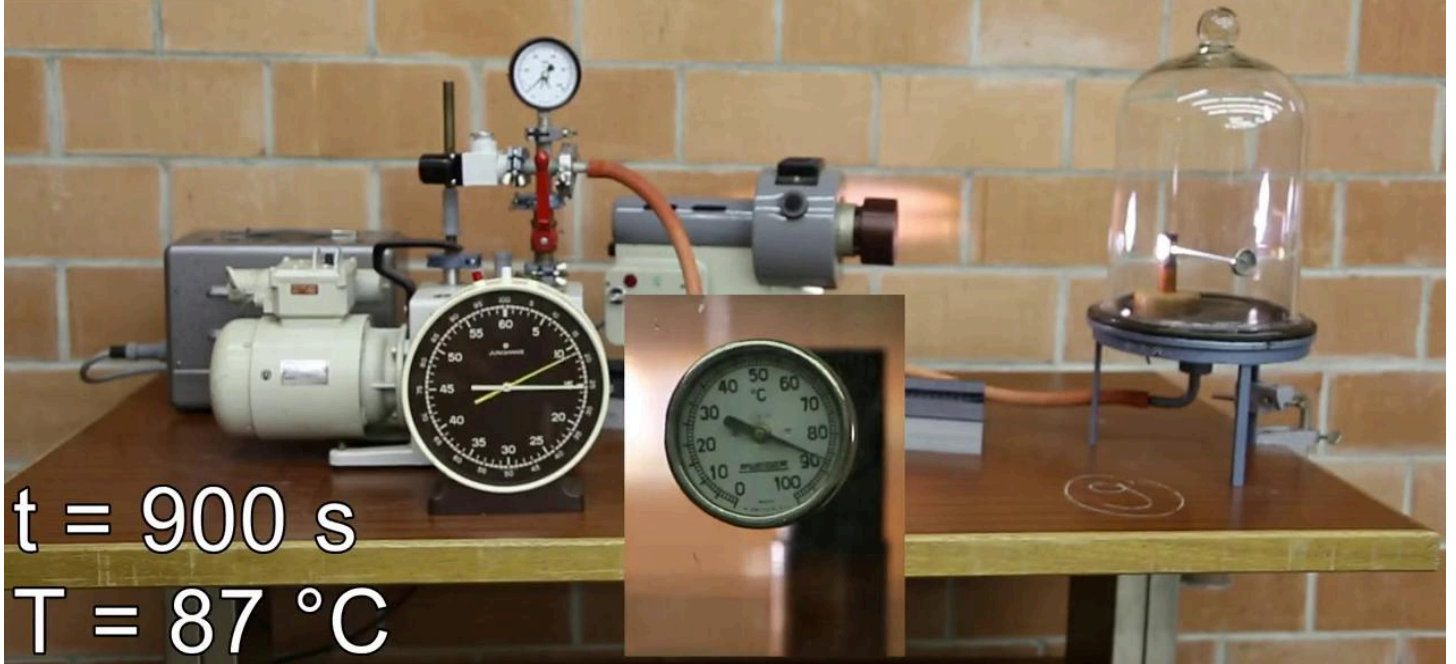
That's the formula you see right now. Therefore, if a material has an influence, it is very small because it has a small specific heat and a small thermal conductivity. It can be very hot, but given to an interface temperature that is not very high. The other way around, if the material has a very good thermal conductivity and a high specific heat, then the temperature of the interface can be very high, very close to the temperature of the material itself.

Notes

Summary



12m 07s



I end with two experiences in relation to the heat transfer by radiation. In this first experience. We will put a thermometer under a vacuum bell. And we're going to light up that thermometer. As there is a vacuum around the thermometer and the support, I ask you to believe it, is a material that conducts heat very badly. We are forced to realize that it is by the radiation that we have a heating of the thermometer. The preparer switches on the pump. And. We start with a room temperature thermometer. Little by little, we observe that the temperature of the thermometer increases.

Notes

Summary





$$\lambda_{max} = \frac{2.9 \cdot 10^{-3} \text{ (m K)}}{T \text{ (K)}}$$

$$\lambda_{max} \approx 600 \text{ nm}$$

$$T \approx 4800 \text{ K}$$

Thermodynamique

We can. Analyze what the position means of this maximum with Wien's law that I recall here. In particular, a maximum was observed at 600 nanometers and with Wien's law, this means that the temperature is 4800 degrees. This should be interpreted as follows. If we had an oven. Equipped with a small hole so that one can observe the radiation in the oven and this oven was heated to 4800 degrees. The same spectrum would be observed that the spectrum of this lamp, as far as we can admit that the spectrum of this lamp has the spectrum of the so-called black body radiation.

Notes

Summary



15m 33s



- Conduction thermique
- Refroidissement avec/
sans ventilation
- Déphasage thermique
- Effusivité
- Rayonnement

Thermodynamique

I summarize. We have seen some experiments concerning the thermal transfer. We have seen a steady state heat transfer. This led us to ask the question heat loss at the surface of a metal with or without forced air. We then illustrated the concept of thermal phase shift by measuring temperature versus time on a bar with heated at one end. Then I showed an experiment of contact between two objects that can be described with the notion of diffusivity. And finally, we observed. A thermal transfer by radiation. Thank you for your attention.

Notes

Summary



16m 23s